

DAFTAR PUSTAKA

- Aqsa Brilianza, Misbahul Jannah, & Abd Mujahid Hamdan. (2020). *Lubang Hitam (Sebuah Pengantar Populer)*. CV. Pustaka Learning Center.
- Fathi, M. (2023). Analytical Study of Particle Geodesics Around a Scale-Dependent De Sitter Black Hole. *Annals of Physics*, 457, 169401. <https://doi.org/10.1016/j.aop.2023.169401>
- Gonzalez Montoya, F. (2024). Impenetrable Barriers in the Phase Space of a Particle Moving Around a Kerr Rotating Black Hole. *Physica D: Nonlinear Phenomena*, 468, 134290. <https://doi.org/10.1016/j.physd.2024.134290>
- Hasanuddin. (2012). Pembelokan Cahaya Dalam Ruang Waktu Kerr-Einsten. *POSITRON*, 2(2), 21–24.
- Hod, S. (2013). Spherical Null Geodesics of Rotating Kerr Black Holes. *Physics Letters B*, 718(4–5), 1552–1556. <https://doi.org/10.1016/j.physletb.2012.12.047>
- J. M. Bardeen, B. Carter, & S. W. Hawking. (1973). The Four Laws of Black Hole Mechanics. *Communications in Mathematical Physics*, 161–170.
- Kesden, M., Sperhake, U., & Berti, E. (2010). Relativistic Suppression of Black Hole Recoils. *The Astrophysical Journal*, 715(2), 1006–1011. <https://doi.org/10.1088/0004-637X/715/2/1006>
- Kunz, J. (2014). Black Holes in Higher Dimensions. *Nuclear Physics B - Proceedings Supplements*, 251–252, 27–32. <https://doi.org/10.1016/j.nuclphysbps.2014.04.005>
- M. P. Hobson, G. P. Efstathiou, & A. N. Lasenby. (2006). *General Relativity: An Introduction for Physicists*. Cambridge University Press.
- Matt Visser. (2008). *The Kerr spacetime: A brief introduction*.
- McInnes, B. (2012). Kerr Black Holes are Not Fragile. *Nuclear Physics B*, 857(3), 362–379. <https://doi.org/10.1016/j.nuclphysb.2011.12.015>
- Moshe Carmeli. (1982). *Classical fields: General relativity and Gauge Theory*. New York: J. Wiley.
- S. Chandra Sekhar. (1983). *The Mathematical Theory of Black Holes*. Oxford University Press.
- Seetharaman, R., Gayathri, S., Tharun, M., Sreeja Mole, S. S., & Anandan, K. (2022). A Study on Black Holes and its Rotation. *Materials Today: Proceedings*, 51, 2380–2383. <https://doi.org/10.1016/j.matpr.2021.11.584>
- Sunkar E. Gautama. (2018). *Pengantar Teori Relativitas Umum dan Kosmologi*. Paradoks Softbook Publisher.

Zhang, M., & Liu, W.-B. (2019). Innermost Stable Circular Orbits of Charged Spinning Test Particles. *Physics Letters B*, 789, 393–398.
<https://doi.org/10.1016/j.physletb.2018.12.051>

LAMPIRAN

A. Perhitungan Persamaan Geodesik

Berangkat dari persamaan (5):

$$I = \int \mathcal{L} ds \quad (\text{A.1})$$

Kemudian divariasikan:

$$\delta I = \delta \int \mathcal{L} ds = \int \delta \mathcal{L} ds \quad (\text{A.2})$$

Mengingat bahwa $\mathcal{L} = \left(g_{\mu\nu} \frac{dx^\mu}{ds} \frac{dx^\nu}{ds} \right)^{\frac{1}{2}}$, maka variasi dari $\mathcal{L}$:

$$\begin{aligned} \delta \mathcal{L} &= \frac{\partial \mathcal{L}}{\partial x^\mu} \delta x^\mu + \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \delta \left(\frac{dx^\mu}{ds} \right) ; \delta \left(\frac{dx^\mu}{ds} \right) = \frac{d}{ds} \delta x^\mu \\ &= \frac{\partial \mathcal{L}}{\partial x^\mu} \delta x^\mu + \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \frac{d}{ds} \delta x^\mu \\ &= \frac{\partial \mathcal{L}}{\partial x^\mu} \delta x^\mu + \frac{d}{ds} \left[\frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \delta x^\mu \right] - \left[\frac{d}{ds} \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \right] \delta x^\mu ds \end{aligned} \quad (\text{A.3})$$

Variasi δx^μ pada batas titik awal dan titik akhir adalah nol, maka suku kedua pada ruas kanan di atas menjadi lenyap.

$$\delta \int \mathcal{L} ds = \int \left[\frac{\partial \mathcal{L}}{\partial x^\mu} - \frac{d}{ds} \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \right] \delta x^\mu ds = 0 \quad (\text{A.4})$$

Karena $\delta x^\mu ds$ bernilai apa saja, maka persamaan geodesik ditentukan oleh:

$$\frac{\partial \mathcal{L}}{\partial x^\mu} - \left[\frac{d}{ds} \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} \right] = 0 \quad (\text{A.5})$$

Sekarang selesaikan bentuk $\frac{\partial \mathcal{L}}{\partial x^\mu}$ dan $\frac{d}{ds} \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)}$, dengan memasukkan nilai $\mathcal{L}$ dan $ds^2 = (g_{\rho\sigma} dx^\rho dx^\sigma)$, diperoleh hubungan:

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial x^\mu} &= \frac{1}{2} \left(g_{\rho\sigma} \frac{dx^\rho}{ds} \frac{dx^\sigma}{ds} \right)^{-\frac{1}{2}} \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} \\ &= \frac{1}{2} \frac{ds}{ds} \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} \\ &= \frac{1}{2} \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} \\ \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds} \right)} &= \left(g_{\rho\sigma} \frac{dx^\rho}{ds} \frac{dx^\sigma}{ds} \right)^{-\frac{1}{2}} g_{\mu\alpha} \frac{dx^\alpha}{ds} \\ &= \frac{ds}{ds} g_{\mu\alpha} \frac{dx^\alpha}{ds} = g_{\mu\alpha} \frac{dx^\alpha}{ds} \end{aligned}$$

$$\begin{aligned}\frac{d}{ds} \frac{\partial \mathcal{L}}{\partial \left(\frac{dx^\mu}{ds}\right)} &= \frac{ds}{ds} \left(g_{\mu\alpha} \frac{dx^\alpha}{ds} \right) \\ &= g_{\mu\alpha} \frac{d^2 x^\alpha}{ds^2} + \frac{\partial g_{\mu\alpha}}{\partial x^\beta} \frac{dx^\alpha}{ds} \frac{ds^\beta}{ds}\end{aligned}$$

Substitusi hasil ke persamaan (A.5) diperoleh:

$$\begin{aligned}\frac{1}{2} \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} - g_{\mu\alpha} \frac{d^2 x^\alpha}{ds^2} - \frac{\partial g_{\mu\alpha}}{\partial x^\beta} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} &= 0 \\ g_{\mu\alpha} \frac{d^2 x^\alpha}{ds^2} + \frac{1}{2} \left(2 \frac{\partial g_{\mu\alpha}}{\partial x^\beta} - \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \right) \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} &= 0 \\ g_{\mu\alpha} \frac{d^2 x^\alpha}{ds^2} + \frac{1}{2} \left(\frac{\partial g_{\mu\alpha}}{\partial x^\beta} + \frac{\partial g_{\mu\beta}}{\partial x^\alpha} - \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \right) \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} &= 0\end{aligned}$$

simbol Christoffel jenis pertama:

$$\Gamma_{\mu\alpha\beta} = \frac{1}{2} \left(\frac{\partial g_{\mu\alpha}}{\partial x^\beta} + \frac{\partial g_{\mu\beta}}{\partial x^\alpha} - \frac{\partial g_{\alpha\beta}}{\partial x^\mu} \right) \quad (\text{A.6})$$

Sehingga persamaan geodesik dapat dituliskan sebagai berikut:

$$\begin{aligned}g_{\mu\nu} \frac{d^2 x^\alpha}{ds^2} + \Gamma_{\alpha\beta\mu} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} &= 0 \\ \frac{d^2 x^\gamma}{ds^2} + \Gamma_{\alpha\beta}^\gamma \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} &= 0\end{aligned} \quad (\text{A.7})$$

B. Perhitungan Geodesik dalam Bidang Ekuator

1. Metrik Kerr dalam Koordinat Boyer-Linquist di Ekuator

Berangkat dari persamaan (3) dimana

$$\begin{aligned}ds^2 &= \frac{-\Delta - a^2 \sin^2 \theta}{\rho^2} dt^2 - 2 \frac{2Mar \sin^2 \theta}{\rho^2} dt d\phi + \frac{\rho^2}{\Delta} dr^2 + \rho^2 d\theta^2 \\ &\quad + \frac{(r^2 + a^2) - a^2 \Delta \sin^2 \theta}{\rho^2} \sin^2 \theta d\phi^2\end{aligned} \quad (\text{B.1})$$

Ditinjau dari ekuator $\theta = \frac{\pi}{2}$, $\Delta = r^2 - 2Mr + a^2$ dan $\rho^2 = r^2 + a^2 \cos^2 \theta$

$$\begin{aligned}ds^2 &= \frac{-[(r^2 - 2Mr + a^2) - a^2]}{r^2} dt^2 - 2 \frac{2Ma}{r} dt d\phi + \frac{r^2}{r^2 - 2Mr + a^2} dr^2 + r^2 d\theta^2 + \\ &\quad \frac{(r^2 + a^2)^2 - a^2(r^2 - 2Mr + a^2)}{r^2} d\phi^2 \\ &= 1 \left(1 - \frac{2M}{r} \right) dt^2 - 2 \frac{2Ma}{r} dt d\phi + \frac{r^2}{\Delta} dr^2 + \frac{(r^2 + a^2)^2 - a^2(r^2 - 2Mr + a^2)}{r^2} d\phi^2\end{aligned}$$

Uraikan $(r^2 + a^2)^2$ dan $a^2(r^2 - 2Mr + a^2)$:

$$(r^2 + a^2)^2 = r^4 + a^2 r^2 + a^2 r^2 + a^4 = r^4 + 2a^2 r^2 + a^4$$

$$a^2(r^2 - 2Mr + a^2) = a^2 r^2 - 2Ma^2 r + a^4$$

menghasilkan:

$$ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 - 2\frac{2Ma}{r} dt d\phi + \frac{r^2}{\Delta} dr^2 + \left[(r^2 + a^2) + \frac{2Ma^2}{r}\right] d\phi^2 \quad (\text{B.2})$$

2. Mencari Komponen Metrik Kontravarian

Komponen metrik kontravarian didapatkan dari menghitung invers metrik kovarian:

$$g^{\mu\nu} = \frac{1}{\det(g_{\mu\nu})} \text{adj}(g_{\mu\nu}) \quad (\text{B.3})$$

Karena elemen diagonal tidak bergantung pada t dan ϕ maka:

$$g_{\mu\nu} = \begin{pmatrix} g_{tt} & g_{t\phi} \\ g_{t\phi} & g_{\phi\phi} \end{pmatrix} \rightarrow g^{\mu\nu} = \frac{1}{|g_{\mu\nu}|} \begin{pmatrix} g_{\phi\phi} & -g_{t\phi} \\ -g_{t\phi} & g_{tt} \end{pmatrix} \quad (\text{B.4})$$

Komponen metrik kovarian didapatkan dari persamaan (B.2) dengan meninjau bentuk umum metrik aksial stationer (1). Menghitung determinan $g_{\mu\nu}$:

$$\begin{aligned} |g_{\mu\nu}| &= g_{tt}g_{\phi\phi} - (g_{t\phi})^2 \\ &= -\left(1 - \frac{2M}{r}\right) \left[(r^2 + a^2) + \frac{2Ma^2}{r}\right] - \left(-\frac{2Ma}{r}\right)^2 \\ &= -\left[\frac{2Ma^2}{r} + (r^2 + a^2) - \frac{2M}{r}(r^2 + a^2) - \frac{4M^2a^2}{r^2}\right] - \frac{4M^2a^2}{r^2} \\ &= -\frac{2Ma^2}{r} - (r^2 - a^2) + 2Mr + \frac{2Ma^2}{r} + \frac{4M^2a^2}{r^2} - \frac{4M^2a^2}{r^2} \\ &= -(r^2 + a^2 - 2Mr) = -\Delta \end{aligned}$$

Maka komponen metrik kontravarian $g^{\mu\nu}$:

$$g^{tt} = \frac{1}{|g_{\mu\nu}|} (g_{\phi\phi}) = -\frac{1}{\Delta} \left[(r^2 + a^2) + \frac{2Ma^2}{r}\right] \quad (\text{B.5})$$

$$g^{t\phi} = \frac{1}{|g_{\mu\nu}|} (-g_{t\phi}) = -\frac{1}{\Delta} \left(-\frac{2Ma}{r}\right) = \frac{2Ma}{r^2} \quad (\text{B.6})$$

$$g^{\phi\phi} = \frac{1}{|g_{\mu\nu}|} (g_{tt}) = -\frac{1}{\Delta} \left[-\left(1 - \frac{2M}{r}\right)\right] = \frac{1}{\Delta} \left(1 - \frac{2M}{r}\right) \quad (\text{B.7})$$

$$g^{rr} = \frac{1}{g_{rr}} = \frac{\Delta}{r^2} \quad (\text{B.8})$$

$$g^{\theta\theta} = \frac{1}{g_{\theta\theta}} = \frac{1}{r^2} \quad (\text{B.9})$$

3. Mencari Komponen Kecepatan Koordinat $\dot{t}$, $\dot{\phi}$, dan $\dot{r}$

Komponen kecepatan koordinat $\dot{t}$, $\dot{\phi}$, dan $\dot{r}$ dalam geodesik partikel di sekitar lubang hitam Kerr dapat dihitung menggunakan persamaan Euler-Lagrange dari Lagrangian relativistik:

$$\mathcal{L} = \frac{1}{2} g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu \quad (\text{B.10})$$

Dengan persamaan Lagrange:

$$\frac{d}{d\lambda} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}^\mu} \right) - \frac{\partial \mathcal{L}}{\partial x^\mu} = 0 \quad (\text{B.11})$$

Untuk E :

$$\begin{aligned} \frac{\partial}{\partial \lambda} \left(\frac{\partial \mathcal{L}}{\partial \dot{t}} \right) - \frac{\partial \mathcal{L}}{\partial t} &= 0 \\ \frac{\partial}{\partial \lambda} \left(\frac{\partial \mathcal{L}}{\partial \dot{t}} \right) &= \frac{\partial \mathcal{L}}{\partial t} \\ \frac{\partial}{\partial \lambda} \left(\frac{\partial}{\partial \dot{t}} \frac{1}{2} (g_{tt} \dot{t}^2 + 2g_{t\phi} \dot{t} \dot{\phi} + g_{rr} \dot{r}^2 + g_{\theta\theta} \dot{\theta}^2 + g_{\phi\phi} \dot{\phi}^2) \right) &= 0 \\ \frac{\partial}{\partial \lambda} (g_{tt} \dot{t} + g_{t\phi} \dot{\phi}) &= 0 \\ g_{tt} \dot{t} + g_{t\phi} \dot{\phi} &= -E \\ -\left(1 - \frac{2M}{r}\right) \dot{t} + \left(-\frac{2Ma}{r}\right) \dot{\phi} &= -E \end{aligned} \quad (\text{B.12})$$

Untuk J :

$$\begin{aligned} \frac{\partial}{\partial \lambda} \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) - \frac{\partial \mathcal{L}}{\partial \phi} &= 0 \\ \frac{\partial}{\partial \lambda} \left(\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right) &= \frac{\partial \mathcal{L}}{\partial \phi} \\ \frac{\partial}{\partial \lambda} \left(\frac{\partial}{\partial \dot{\phi}} \frac{1}{2} (g_{tt} \dot{t}^2 + 2g_{t\phi} \dot{t} \dot{\phi} + g_{rr} \dot{r}^2 + g_{\theta\theta} \dot{\theta}^2 + g_{\phi\phi} \dot{\phi}^2) \right) &= 0 \\ \frac{\partial}{\partial \lambda} (g_{\phi\phi} \dot{\phi} + g_{t\phi} \dot{t}) &= 0 \\ g_{\phi\phi} \dot{\phi} + g_{t\phi} \dot{t} &= J \\ \left[(r^2 + a^2) + \frac{2Ma^2}{r} \right] \dot{\phi} + \left(-\frac{2Ma}{r} \right) \dot{t} &= J \end{aligned} \quad (\text{B.13})$$

Kemudian eliminasi untuk $\dot{t}$ dan $\dot{\phi}$

Komponen $\dot{t}$:

$$-\left(1 - \frac{2M}{r}\right)\dot{t} + \left(-\frac{2Ma}{r}\right)\dot{\phi} = -E$$

$$\left(-\frac{2Ma}{r}\right)\dot{\phi} = -E + \left(1 - \frac{2M}{r}\right)\dot{t}$$

$$\dot{\phi} = \left[-E + \left(1 - \frac{2M}{r}\right)\dot{t}\right]\left(-\frac{r}{2Ma}\right)$$

Substitusi nilai $\dot{\phi}$ ke persamaan (B.13)

$$\begin{aligned} & \left[(r^2 + a^2) + \frac{2Ma^2}{r}\right]\dot{\phi} + \left(-\frac{2Ma}{r}\right)\dot{t} = J \\ & \left[(r^2 + a^2) + \frac{2Ma^2}{r}\right]\left[-E + \left(1 - \frac{2M}{r}\right)\dot{t}\right]\left(-\frac{r}{2Ma}\right) + \left(-\frac{2Ma}{r}\right)\dot{t} = J \\ & -\frac{2Ma}{r}\dot{t} - \frac{r}{2Ma}\left[-E(r^2 + a^2) - \frac{2Ma^2}{r}E + (r^2 + a^2)\left(1 - \frac{2M}{r}\right)\dot{t} \right. \\ & \quad \left. + \frac{2Ma^2}{r}\left(1 - \frac{2M}{r}\right)\dot{t}\right] = J \\ & -\frac{2Ma}{r}\dot{t} + \frac{r(r^2 + a^2)E}{2Ma} + aE - \frac{r}{2Ma}(r^2 + a^2)\left(1 - \frac{2M}{r}\right)\dot{t} - a\left(1 - \frac{2M}{r}\right)\dot{t} = J \\ & \dot{t}\left[-\frac{2Ma}{r} - \frac{r}{2Ma}(r^2 + a^2)\left(1 - \frac{2M}{r}\right) - a\left(1 - \frac{2M}{r}\right)\right] = J - \frac{r(r^2 + a^2)E}{2Ma} - aE \\ & \dot{t}\left(-\frac{r}{2Ma}\Delta\right) = J - aE - \frac{r(r^2 + a^2)E}{2Ma} \\ & \dot{t} = -\frac{1}{\Delta}\left(\frac{2Ma}{r}J - \frac{2Ma^2E}{r} - (r^2 + a^2)E\right) \\ & \dot{t} = \frac{1}{\Delta}\left[\left((r^2 + a^2) + \frac{2Ma^2}{r}\right)E - \frac{2Ma}{r}J\right] \end{aligned} \tag{B.14}$$

Komponen $\dot{\phi}$:

$$-\left(1 - \frac{2M}{r}\right)\dot{t} + \left(-\frac{2Ma}{r}\right)\dot{\phi} = -E$$

$$-\left(1 - \frac{2M}{r}\right)\dot{t} = -E + \left(\frac{2Ma}{r}\right)\dot{\phi}$$

$$\dot{t} = \left[-E + \left(\frac{2Ma}{r}\right)\dot{\phi}\right]\left(-1 + \frac{r}{2M}\right)$$

$$\dot{t} = \left[-E + \left(\frac{2Ma}{r} \right) \dot{\phi} \right] \left(-\frac{r}{r-2M} \right)$$

Substitusi nilai $\dot{t}$ ke persamaan (B.12)

$$\begin{aligned} & \left[(r^2 + a^2) + \frac{2Ma^2}{r} \right] \dot{\phi} + \left(-\frac{2Ma}{r} \right) \dot{t} = J \\ & \left[(r^2 + a^2) + \frac{2Ma^2}{r} \right] \dot{\phi} + \left(-\frac{2Ma}{r} \right) \left[-E + \frac{2Ma}{r} \dot{\phi} \right] \left(-\frac{r}{r-2M} \right) = J \\ & \left[(r^2 + a^2) + \frac{2Ma^2}{r} \right] \dot{\phi} + \left[\frac{2MaE}{r} - \frac{4M^2a^2}{r^2} \dot{\phi} \right] \left(-\frac{r}{r-2M} \right) = J \\ & \left[(r^2 + a^2) + \frac{2Ma^2}{r} \right] \dot{\phi} + \left[\left(-\frac{2MaE}{r-2M} \right) + \frac{4M^2a^2}{r(r-2M)} \dot{\phi} \right] = J \\ & \dot{\phi} \left[(r^2 + a^2) + \frac{2Ma^2}{r} + \frac{4M^2a^2}{r(r-2M)} \right] = J + \frac{2MaE}{r-2M} \\ & \dot{\phi} [r(r-2M)(r^2 + a^2) + 2Ma^2(r-2M) + 4M^2a^2] = r(r-2M) \left[J + \frac{2MaE}{r-2M} \right] \\ & \dot{\phi} [r^4 + a^2r^2 - 2Mr^3 - 2Ma^2r + 2Ma^2r - 4M^2a^2 + 4M^2a^2] \\ & \quad = r(r-2M)J + 2MarE \\ & \dot{\phi} [r^4 + a^2r^2 - 2Mr^3] = r(r-2M)J + 2MarE \\ & \dot{\phi} r^2(r^2 + a^2 - 2Mr) = r(r-2M)J + 2MarE \\ & \dot{\phi} r^2\Delta = r(r-2M)J + 2MarE \\ & \dot{\phi} = \frac{1}{\Delta} \left(1 - \frac{2M}{r} \right) J + \frac{2MaE}{r} \end{aligned} \tag{B.15}$$

Komponen koordinat $\dot{r}$:

Dengan menggunakan $g^{\mu\nu}P_\mu P_\nu = \epsilon^2$

$$g^{tt}(P_t)^2 + 2g^{t\phi}P_tP_\phi + g^{\phi\phi}(P_\phi)^2 + g^{rr}(P_r)^2 = \epsilon^2$$

Substitusi $P_t = -E$, $P_t = -E$ dan $P_r = g_{rr} \frac{dr}{d\lambda} = g_{rr}\dot{r}$

$$g^{tt}E^2 - 2g^{t\phi}EJ + g^{\phi\phi}J^2 + g_{rr}\dot{r} = -\epsilon^2$$

$$g_{rr}\dot{r} = -\epsilon^2 - g^{tt}E^2 + 2g^{t\phi}EJ - g^{\phi\phi}J^2$$

$$\dot{r} = \frac{1}{g_{rr}} [-\epsilon^2 - g^{tt} E^2 + 2g^{t\phi} EJ - g^{\phi\phi} J^2]$$

Kemudian substitusi komponen-komponen metrik:

$$\begin{aligned} \dot{r}^2 &= \frac{\Delta}{r^2} \left[\epsilon^2 + \frac{1}{\Delta} \left(r^2 + a^2 + \frac{2Ma^2}{r} \right) E^2 + \frac{4MaEJ}{\Delta r} - \frac{1}{\Delta} \left(1 - \frac{2M}{r} \right) J^2 \right] \\ &= -\frac{\Delta \epsilon^2}{r^2} + \frac{1}{r^2} \left(r^2 + a^2 + \frac{2Ma^2}{r} \right) E^2 + \frac{4MaEJ}{r^3} - \frac{1}{r^2} \left(1 - \frac{2M}{r} \right) J^2 \\ &= -\frac{(r^2 - 2Mr + a^2)\epsilon^2}{r^2} + \frac{E^2}{r^2} \left(r^2 + a^2 + \frac{2Ma^2}{r} \right) + \frac{4MaEJ}{r^3} - \frac{J^2}{r^2} \left(1 - \frac{2M}{r} \right) \\ &= \epsilon^2 + \frac{2M\epsilon^2}{r} - \frac{\epsilon^2 a^2}{r^2} + E^2 + \frac{2Ma^2 E^2}{r^3} - \frac{4MaEJ}{r^2} - \frac{J^2}{r^2} + \frac{2MJ^2}{r^3} \\ &= E^2 - \epsilon^2 + \frac{2M\epsilon^2}{r} + \frac{a^2(-\epsilon^2 + E^2) - J^2}{r^2} + \frac{2M(E^2 a^2 - 2aEJ + J^2)}{r^3} \\ &= E^2 - \epsilon^2 + \frac{2M\epsilon^2}{r} + \frac{a^2(-\epsilon^2 + E^2) - J^2}{r^2} + \frac{2M(aE - J)^2}{r^3} \end{aligned} \quad (B.16)$$

4. Persamaan Lintasan Partikel Bermassa di Ekuator

Dengan meninjau kembali persamaan (B.16), untuk partikel bermassa ($\epsilon^2 = 1$), persamaan menjadi:

$$\dot{r}^2 = (E^2 - 1) + \frac{2M}{r} + \frac{a^2(E^2 - 1) - J^2}{r^2} + \frac{2M(aE - J)^2}{r^3} \quad (B.17)$$

Kemudian menyusun ulang persamaan, diperoleh:

$$\dot{r}^2 - \frac{2M}{r} - \frac{a^2(E^2 - 1) - J^2}{r^2} - \frac{2M(aE - J)^2}{r^3} = (E^2 - 1) \quad (B.18)$$

Selanjutnya bentuk persamaan dapat ditulis sebagai:

$$\frac{1}{2} \dot{r}^2 - \frac{M}{r} - \frac{a^2(E^2 - 1)J^2}{2r^2} - \frac{M(aE - J)^2}{r^3} = \frac{1}{2} (E^2 - 1) \quad (B.19)$$

Dengan potensial efektif (V_{eff}):

$$(V_{eff}(r, J, E) = -\frac{M}{r} + \frac{J^2 - a^2(E^2 - 1)}{2r^2} - \frac{M(aE - J)^2}{r^3})$$

C. Gerak Melingkar Partikel Bermassa di Ekuator

Untuk gerak melingkar partikel dipilih $\dot{r} = 0$ dan $\ddot{r} = 0$. Tinjau potensial efektif:

$$V_{eff} = -\frac{M}{r} + \frac{J^2 - a^2(E^2 - 1)}{2r^2} - \frac{M(aE - J)^2}{r^3}$$

Dan

$$\frac{1}{2} \dot{r} + V_{eff} = \frac{1}{2} (E^2 - 1)$$

$$V_{eff} = \frac{1}{2}(E^2 - 1); \quad \frac{dV_{eff}}{dr} \Big|_{r=r_0} = 0$$

Dengan menggunakan $U = \frac{1}{r}$ dan disubstitusi ke persamaan ():

$$-MU + \frac{1}{2}U^2[J^2 - a^2(E^2 - 1)] - M(aE - J)^2U^3 = \frac{1}{2}(E^2 - 1) \quad (C.1)$$

Diferensialkan terhadap U :

$$\begin{aligned} \frac{d}{dU} \left(-MU + \frac{1}{2}U^2[J^2 - a^2(E^2 - 1)] - M(aE - J)^2U^3 \right) &= \frac{d}{dU} \left(\frac{1}{2}(E^2 - 1) \right) \\ -M + U[J^2 - a^2(E^2 - 1)] - 3M(aE - J)^2U^2 &= 0 \end{aligned}$$

Kemudian masukkan variabel baru $x = (aE - J) \Rightarrow J^2 = x^2 + a^2E^2 + 2axE$, maka persamaan (C.1):

$$\begin{aligned} -MU + \frac{1}{2}U^2 \left[(x^2 + a^2E^2 + 2axE) - a^2(E^2 - 1) \right] - M(aE - J)^2U^3 &= \frac{1}{2}(E^2 - 1) \\ -MU + \frac{1}{2}U^2(x^2 + a^2E^2 + 2axE - a^2E^2 + a^2) - Mx^2U^3 &= \frac{1}{2}(E^2 - 1) \\ -MU + \frac{1}{2}(x^2 + 2axE + a^2)U^2 - Mx^2U^3 &= \frac{1}{2}(E^2 - 1) \end{aligned} \quad (C.2)$$

Turunkan persamaan (C.2) terhadap U :

$$\begin{aligned} \frac{d}{du} \left(-MU + \frac{1}{2}U^2(x^2 + a^2E^2 + 2axE - a^2E^2 + a^2) - Mx^2U^3 \right) &= \frac{d}{du} \left(\frac{1}{2}(E^2 - 1) \right) \\ -M + (x^2 + 2axE + a^2)U - 3Mx^2U^2 &= 0 \end{aligned} \quad (C.3)$$

Pers (C.2) masing-masing dikalikan dengan 2:

$$\begin{aligned} -MU + \frac{1}{2}(x^2 + 2axE + a^2)U^2 - Mx^2U^3 &= \frac{1}{2}(E^2 - 1) \\ -2MU + (x^2 + 2axE + a^2)U^2 - 2Mx^2U^3 + 1 &= E^2 \end{aligned} \quad (C.4)$$

Pers (C.3) dikalikan dengan U :

$$\begin{aligned} -MU + (x^2 + 2axE + a^2)U^2 - 3Mx^2U^3 &= 0 \\ (x^2 + 2axE + a^2)U^2 &= MU + 3Mx^2U^3 \end{aligned} \quad (C.5)$$

Substitusi (C.4) ke (C.5):

$$-2MU + MU + 3Mx^2U^3 - 2Mx^2U^3 + 1 = E^2$$

$$\begin{aligned}
 -MU + Mx^2U^3 + 1 &= E^2 \\
 E^2 &= (1 - MU) + Mx^2U^3
 \end{aligned} \tag{C.6}$$

Persamaan (C.6) disubstitusi ke pers (C.4)

$$\begin{aligned}
 -2MU + (x^2 + 2axE + a^2)U^2 - 2Mx^2U^3 + 1 &= (1 - MU) + Mx^2U^3 \\
 -2MU + (x^2 + 2axE + a^2)U^2 - 2Mx^2U^3 + 1 - 1 + MU - Mx^2U^3 &= 0 \\
 -MU + (x^2 + 2axE + a^2)U^2 - 3Mx^2U^3 &= 0 \\
 -M + (x^2 + 2axE + a^2)U - 3Mx^2U^2 &= 0 \\
 -M + x^2U + 2axEU + a^2U - 3Mx^2U^2 &= 0 \\
 2axEU &= 3Mx^2U^2 + M - x^2U - a^2U \\
 2axEU &= x^2U(3MU - 1) - (a^2 - M)
 \end{aligned} \tag{C.7}$$

Dengan mengeliminasi E di antara persamaan (C.6) dan (C.7) didapatkan:

$$\begin{aligned}
 4a^2x^2U^2[(1 - MU) + Mx^2U^3] &= x^4U^2(3MU - 1)^2 + (a^2U - M)^2 \\
 &\quad - 2x^2U(3MU - 1)(a^2U - M) \\
 4a^2x^2U^2(1 - MU) + 4Ma^2x^4U^5 &= x^4U^2(3MU - 1)^2 + (a^2U - M)^2 \\
 &\quad - 2x^2U(3MU - 1)(a^2U - M) \\
 x^4[(3MU - 1)^2 - 4Ma^2U^3]U^2 + x^2[4a^2U(1 - MU) - 2(3MU - 1)(a^2U - M)]U + (a^2U - M)^2 &= 0
 \end{aligned}$$

atau

$$U^2[(3MU - 1)^2 - 4Ma^2U^3]x^4 + 2U[2a^2U(MU - 1) - (3MU - 1)(a^2U - M)]x^2 + (a^2U - M)^2 = 0$$

atau misalkan $Z = x^2$

$$\begin{aligned}
 &U^2[(3MU - 1)^2 - 4Ma^2U^3]Z^2 + \\
 &2U[2a^2U(MU - 1) - (3MU - 1)(a^2U - M)]Z + (a^2U - M)^2 = 0
 \end{aligned} \tag{C.8}$$

Gunakan rumus kuadrat:

$$Z_{12} = \frac{-B \pm \sqrt{B^2 - 4AC}}{2A} \tag{C.9}$$

Pertama hitung nilai A, B , dan C :

$$A = U^2[(3MU - 1)^2 - 4Ma^2U^3]$$

$$= U^2(9M^2U^2 - 6MU + 1 - 4Ma^2U^3)$$

$$B = 2U[2a^2U(MU - 1) - (3MU - 1)(a^2U - M)]$$

$$= 2U(2a^2MU^2 - 2a^2U - 3a^2MU^2 + 3M^2U + a^2U - M)$$

$$= 2U(3M^2U - a^2MU^2 - a^2U - M)$$

$$= 6M^2U^2 - 2a^2MU^3 - 2a^2U^2 - 2MU$$

$$C = (a^2U - M)^2$$

$$= (a^2U - M)(a^2U - M)$$

$$= a^4U^2 - a^2MU - a^2MU + M$$

$$= a^4U^2 - 2a^2MU + M$$

Substitusi nilai A, B, dan C ke dalam rumus kuadrat, didapatkan:

$$Z_{12} = \frac{(a\sqrt{U} \pm \sqrt{M})^2}{U(1-3MU \mp 2a\sqrt{MU^3})} \quad (\text{C.10})$$

Mengingat $Z = x^2$, maka Jarak radial x dalam koordinat Boyer-Lindquist dinyatakan sebagai:

$$x^2 = \frac{(a\sqrt{U} \pm \sqrt{M})^2}{U(1-3MU \mp 2a\sqrt{MU^3})}$$

$$x = \frac{\sqrt{(a\sqrt{U} \pm \sqrt{M})^2}}{\sqrt{U(1-3MU \mp 2a\sqrt{MU^3})}U} = - \frac{a\sqrt{U} \pm \sqrt{M}}{[U(1-3MU \mp 2a\sqrt{MU^3})]^{\frac{1}{2}}} \quad (\text{C.11})$$

Kemudian solusi di atas disubstitusi ke dalam persamaan (39)

$$E^2 = (1 - MU) + Mx^2U^3$$

$$= (1 - MU) + M \left(- \frac{(a\sqrt{U} \pm \sqrt{M})^2}{U(1-3MU \mp 2a\sqrt{MU^3})} \right) U^3$$

$$= (1 - MU) + M \left(- \frac{a^2U + M \pm 2a\sqrt{UM}}{U(1-3MU \mp 2a\sqrt{MU^3})} \right) U^3$$

$$= \frac{(1-MU)(1-3MU \mp 2a\sqrt{MU^3}) + MU^2(a^2U + M \pm 2a\sqrt{UM})}{(1-3MU \mp 2a\sqrt{MU^3})}$$

$$= \frac{1-3MU \mp 2a\sqrt{MU^3} - MU + 3M^2U^2 \pm 2aMU\sqrt{MU^3} + MU^2a^2 + M^2U^2 \pm 2aMU\sqrt{U^3M}}{1-3MU \mp 2a\sqrt{MU^3}}$$

Suku pada pembilang dapat disederhanakan:

$$1 - 4MU \mp 2a\sqrt{MU^3} + 4M^2U^2 \pm 4aMU\sqrt{MU^3} + MU^3a^2$$

Karena $(a + b + c)^2 = a^2 + b^2 + c^2 + 2ab + 2ac + 2bc$ maka diperoleh:

$$= (-1 + 2MU \pm a\sqrt{MU^3})^2$$

Kemudian masukkan kembali ke persamaan awal:

$$\begin{aligned} E^2 &= \frac{(-1+2MU \pm a\sqrt{MU^3})^2}{(1-3MU \mp 2a\sqrt{MU^3})} \\ E &= \frac{(-1+MU \pm a\sqrt{MU^3})}{(1-3MU \mp 2a\sqrt{MU^3})^{\frac{1}{2}}} \end{aligned} \quad (C.12)$$

Selanjutnya untuk nilai J :

$$x = aE - J$$

$$J = x + aE$$

$$\begin{aligned} J &= \left(-\frac{a\sqrt{U} \pm \sqrt{M}}{\left[U(1-3MU \mp 2a\sqrt{MU^3}) \right]^{\frac{1}{2}}} \right) + a \left(\frac{(-1+MU \pm a\sqrt{MU^3})}{(1-3MU \mp 2a\sqrt{MU^3})^{\frac{1}{2}}} \right) \\ J &= \frac{-a\sqrt{U} \pm \sqrt{M} + a\sqrt{U}(-1+MU \pm a\sqrt{MU^3})}{\sqrt{U}(1-3MU \mp 2a\sqrt{MU^3})^{\frac{1}{2}}} \\ J &= \mp \frac{\sqrt{M}(1+a^2U^2 \pm 2a\sqrt{MU^3})}{\sqrt{U}(1-3MU \mp 2a\sqrt{MU^3})^{\frac{1}{2}}} \end{aligned} \quad (C.13)$$

D. Persamaan ISCO

Dari persamaan (50):

$$U = \frac{J^2 + a^2(E^2 - 1)}{6Mx^2}$$

Substitusi nilai J^2 , E^2 dan x^2 dari persamaan (44, 45 dan 46):

Untuk x_1, J_1, E_1 , Misal $Q_- = 1 - 3MU - 2a\sqrt{MU^3}$

$$J^2 = \frac{M(1+a^2U^2 \pm 2a\sqrt{MU^3})^2}{U(1-3MU - 2a\sqrt{MU^3})} = \frac{M(1+a^2U^2 \pm 2a\sqrt{MU^3})^2}{UQ_-} \quad (D.1)$$

$$E^2 = \frac{(1-2MU+a\sqrt{MU^3})^2}{1-3MU-2a\sqrt{MU^3}} = \frac{(1-2MU+a\sqrt{MU^3})^2}{Q_-} \quad (D.2)$$

$$x^2 = \frac{(a\sqrt{U}+\sqrt{M})^2}{U(1-3MU-2a\sqrt{MU^3})} = \frac{a^2U+M+2a\sqrt{MU}}{UQ_-} \quad (D.3)$$

$$U = \left[\frac{M(1+a^2U^2+2a\sqrt{MU^3})^2}{UQ_-} + a^2 \left(\frac{(1-2MU+a\sqrt{MU^3})^2}{Q_-} - 1 \right) \right] \times \frac{UQ_-}{6M(a^2U+M+2a\sqrt{MU})}$$

$$\begin{aligned} 6MU(a^2U+M+ &= M(1+a^2U^2+2a\sqrt{MU^3})^2 - a^2U(1-2MU+ \\ \sqrt{MU}) & a\sqrt{MU^3})^2 + a^2U(1-3MU-2a\sqrt{MU^3}) \\ &= M(1+a^4U^4+4Ma^2U^3+2a^2U^2+4a\sqrt{MU^3}+ \\ & 4a^3U^2\sqrt{MU^3}) - a^2U(1+4M^2U^2+Ma^2U^3-4MU- \\ & 2a\sqrt{MU^3}+4MaU\sqrt{MU^3}) + a^2U-3Ma^2U^2- \\ & 2a^3U\sqrt{MU^3} \\ &= M+Ma^4U^4+4M^2a^2U^2+2Ma^2U^2+4Ma\sqrt{MU^3}+ \\ & 4Ma^3U^2\sqrt{MU^3}-a^2U-4M^2a^2U^3-Ma^4U^4+ \\ & 4Ma^2U^2+2a^3U\sqrt{MU^3}-4Ma^2U^2\sqrt{MU^3}+a^2U- \\ & 3Ma^2U^2-2a^3U\sqrt{MU^3} \\ &= M+3Ma^2U^2+4Ma\sqrt{MU^3}+8Ma^3U^2\sqrt{MU^3}- \\ & 2a^3U\sqrt{MU^3} \end{aligned}$$

$$3Ma^2U^3+6M^2U+8MaU\sqrt{MU}-M=0$$

atau

$$3a^2U^2+6MU+8a\sqrt{MU^3}-1=0$$

Masukkan $U = \frac{1}{r}$:

$$r^2-6Mr-3a^2-8a\sqrt{Mr}=0 \quad (D.4)$$

Dengan cara yang sama untuk x_2, J_2, E_2 misal $Q_+ = 1-3MU+2a\sqrt{MU^3}$

$$x^2 = \frac{(a\sqrt{U}-\sqrt{M})^2}{U(1-3MU+2a\sqrt{MU^3})} = \frac{a^2U+M-2a\sqrt{MU}}{UQ_+} \quad (D.5)$$

$$E^2 = \frac{(1-2MU+a\sqrt{MU^3})^2}{1-3MU+2a\sqrt{MU^3}} = \frac{(1-2MU+a\sqrt{MU^3})^2}{Q_+} \quad (D.6)$$

$$J^2 = \frac{M(1+a^2U^2-2a\sqrt{MU^3})^2}{U(1-3MU+2a\sqrt{MU^3})} = \frac{M(1+a^2U^2-2a\sqrt{MU^3})^2}{UQ_+} \quad (D.7)$$

Maka:

$$r^2-6Mr-3a^2+8a\sqrt{Mr}=0 \quad (D.8)$$

Dengan demikian persamaan ISCO didapatkan:

$$r^2 - 6Mr - 3a^2 \pm 8a\sqrt{Mr} = 0$$